## Reciprocal Frames, the Vector Derivative and Curvilinear Coordinates.

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## Reciprocal Frames



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We call such a frame a reciprocal frame. Note that since any vector $a$ can be written as $a=a^{k} e_{k} \equiv \sum a^{k} e_{k}$ (ie we are adopting the convention that repeated indices are summed over), we have

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Thus we would be able to recover the components of a given vector in a similar way to that used for orthonormal frames.
So how do we find a reciprocal frame?

## Reciprocal Frames cont....

We need, for example, $e^{1}$ to be orthogonal to the set of vectors $\left\{e_{2}, e_{3}, \ldots, e_{n}\right\}$. ie $e^{1}$ must be perpendicular to the hyperplane $e_{2} \wedge e_{3} \wedge \ldots . \wedge e_{n}$.

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We find this by dualisation, ie multiplication by $I$ [note: $I$ is the $n$-d pseudoscalar for our space]. We form $e^{1}$ via

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$\alpha$ is a scalar found by dotting with $e_{1}$ :

$$
e_{1} \cdot e^{1}=1=e_{1} \cdot\left(\alpha e_{2} \wedge e_{3} \wedge \ldots \wedge e_{n} I\right)=\alpha\left(e_{1} \wedge e_{2} \wedge \ldots \wedge e_{n}\right) I
$$

(this uses a useful GA relation $a \cdot(B I)=(a \wedge B) I)$.

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If we let

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we see that $\alpha E_{n} I=1$, so that $\alpha=E_{n}^{-1} I^{-1}$. Thus giving us

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These reciprocal frames are remarkably useful!

## Exercises 1

(1) Show that $a \cdot(B I)=(a \wedge B) I$. [Hint: make use of the fact that $\left.a \cdot\left(B_{r} I_{n}\right)=\left\langle a B_{r} I_{n}\right\rangle_{n-r-1}\right]$.
(2) For $\left\{f_{1}, f_{2}, f_{3}\right\}=\left\{e_{1}, e_{1}+2 e_{3}, e_{1}+e_{2}+e_{3}\right\}$ show, using the given formulae, that the reciprocal frame is given by

$$
\left\{f^{1}, f^{2}, f^{3}\right\}=\left\{e_{1}-\frac{1}{2}\left(e_{2}+e_{3}\right), \frac{1}{2}\left(e_{3}-e_{2}\right), e_{2}\right\}
$$

[these are the reciprocal frames shown in the earlier pictures]

## Exercises 2

(1) Interchanging the role of frame and reciprocal frame, verify that we can write the frame vectors as

$$
\begin{aligned}
& \quad e_{k}=(-1)^{k+1} e^{1} \wedge e^{2} \wedge \ldots \wedge e^{k} \wedge \ldots \wedge e^{n}\left\{E^{n}\right\}^{-1} \\
& \text { where } E^{n}=e^{1} \wedge e^{2} \wedge \ldots \wedge e^{n} \neq 0 .
\end{aligned}
$$

(2) Now show that we can move vectors through each other (changing sign) to give

$$
e_{k}=(-1)^{k-1} e^{n} \wedge e^{n-1} \wedge \ldots \wedge \check{e}^{k} \wedge \ldots \wedge e^{1}\{I V\}
$$

where $\left\{E^{n}\right\}^{-1}=I V$, and $V$ is therefore a volume factor.

## Example: Recovering a Rotor in 3-d

As an example of using reciprocal frames, consider the problem of recovering the rotor which rotates between two 3-d non-orthonormal frames $\left\{e_{k}\right\}$ and $\left\{f_{k}\right\}$, ie find $R$ such that

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A very easy way of recovering rotations.

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A vector $x$ can be represented in terms of coordinates in two ways:

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(Summation implied). Depending on whether we expand in terms of a given frame $\left\{e_{k}\right\}$ or its reciprocal $\left\{e^{k}\right\}$. The coefficients in these two frames are therefore given by

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ie the derivative with respect to the first coordinate, keeping the second and third coordinates constant.

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It follows now that if we $\operatorname{dot} \nabla$ with $a$, we get the directional derivative in the $a$ direction:

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We will see later that the definition of $\nabla$ is independent of the choice of frame.

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A Scalar Field $\phi(x)$ : it gives $\nabla \phi$ which is the gradient.

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See later discussions of electromagnetism.

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For cylindrical polars, our position vector is defined in terms of two lengths $\rho, z$ and an angle $\phi: \Longrightarrow$ coordinates are $\left(\rho, \phi_{\underline{1}} z\right) \cdot$

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Thus, we can construct a second, reciprocal, frame from the coordinates using the vector derivative

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Given coordinates $\left\{x^{i}, i=1, \ldots, n\right\}$, which any position vector, $\mathbf{r}$ [note, use boldface to distinguish from distance from origin], can be expressed in terms of, we can define a set of frame vectors as

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Given coordinates $\left\{x^{i}, i=1, \ldots, n\right\}$, which any position vector, $\mathbf{r}$ [note, use boldface to distinguish from distance from origin], can be expressed in terms of, we can define a set of frame vectors as

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e_{i}(\mathbf{r})=\frac{\partial \mathbf{r}}{\partial x^{i}}
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$$

We see therefore that the vector derivative is crucial in relating coordinates to frames - and we will see how this simplifies manipulations in curvilinear coordinates.

## Div, Grad, Curl in Curvilinear Coordinates

Gradient of a Scalar Function, $\psi$

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\nabla \psi=e^{i} \frac{\partial \psi}{\partial x^{i}} \quad[\text { vector }]
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\nabla \cdot J=e^{i} \frac{\partial}{\partial x^{i}} \cdot\left(J^{j} e_{j}\right)=e^{i} \cdot \frac{\partial\left(J^{j} e_{j}\right)}{\partial x^{i}} \quad[\text { scalar }]
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## Div, Grad, Curl cont....

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Now, take the pseudovector ( $n-1$-blade)
$P=(-1)^{k-1} e^{n} \wedge e^{n-1} \wedge \ldots \wedge e^{k} \wedge \ldots \wedge e^{1}$, and recall that $e_{i}=$ PIV [See Exercises 2]. So that (where $\langle X\rangle$ denotes the scalar part of $X$ )

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$$
\nabla \cdot e_{i}=\langle\nabla(P I V)\rangle=\langle(\nabla P) I V\rangle+\langle P I(\nabla V)\rangle
$$

## Div, Grad, Curl cont....

After some manipulation (which will be outlined in the following exercises) we are able to write

$$
\begin{gathered}
\nabla \cdot J=e_{i} \cdot\left(\nabla J^{i}\right)+J^{i}\left(e_{i} \cdot \nabla(\ln V)\right) \\
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$$
\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i} \quad[\text { bivector }]
$$

## Exercises 3

(1) Since $I V$ is a pseudoscalar, show that

$$
\langle(\nabla P) I V\rangle=\langle(\nabla \wedge P) I V)\rangle
$$

(2) Using the fact that $e_{i}=P I V$, show that

$$
P I(\nabla V)=e_{i} \nabla(\ln V)
$$

(3) Verify that $\nabla \wedge a \wedge b=(\nabla \wedge a) \wedge b-a \wedge(\nabla \wedge b)$, and then, using our previous result of $\nabla \wedge e^{i}=0$, show that

$$
\nabla \wedge P=0
$$

and therefore that

$$
\nabla \cdot e_{i}=e_{i} \cdot \nabla(\ln V)
$$

## Exercises 4

(1) By expanding $\nabla \wedge J$ as

$$
\nabla \wedge J=\nabla \wedge\left(J_{i} e^{i}\right)=\dot{\nabla} \wedge\left(\dot{J}_{i} e^{i}\right)+\dot{\nabla} \wedge\left(J_{i} e^{i}\right)
$$

explain how we obtain the result $\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i}$

## An Example: Spherical Polars in 3d

Recall our coordinates are $(r, \theta, \phi)$, and we also have an orthogonal set of unit vectors ( $\hat{e}_{r}, \hat{e}_{\theta}, \hat{e}_{\phi}$ ) as shown in the diagram. Thus, we can define a frame via $e_{i}=\frac{\partial r}{\partial x^{i}}$ to be

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e_{\phi}=\frac{\partial \mathbf{r}}{\partial \phi}=\frac{\partial\left(r \hat{e}_{r}\right)}{\partial \phi}=r \frac{\partial \hat{e}_{r}}{\partial \phi}=r \frac{\partial\left(\cos \theta \hat{e}_{z}+\sin \theta \hat{e}_{\rho}\right)}{\partial \phi}=r \sin \theta \hat{e}_{\phi}
\end{gathered}
$$

## An Example: Spherical Polars in 3d cont...

From the definition of reciprocal frame we therefore see that the reciprocal vectors are given by

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(check that $\left.e^{i} \cdot e_{j}=\delta_{j}^{i}\right)$.

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Gradient

$$
\nabla \psi=e^{i} \frac{\partial \psi}{\partial x^{i}}=\frac{\partial \psi}{\partial r} \hat{e}_{r}+\frac{1}{r} \frac{\partial \psi}{\partial \theta} \hat{e}_{\theta}+\frac{1}{r \sin \theta} \frac{\partial \psi}{\partial \phi} \hat{e}_{\phi}
$$

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$$

Which agrees with the formula given in tables etc.

## An Example: Spherical Polars in 3d cont...

Divergence

$$
\begin{aligned}
\nabla \cdot J=\frac{1}{V} \frac{\partial\left(V J^{i}\right)}{\partial x^{i}} & =\frac{1}{r^{2} \sin \theta} \frac{\partial\left(r^{2} \sin \theta J^{r}\right)}{\partial r}+\frac{1}{r^{2} \sin \theta} \frac{\partial\left(r^{2} \sin \theta J^{\theta}\right)}{\partial \theta}+\frac{1}{r^{2} \sin \theta} \frac{\partial\left(r^{2} \sin \theta J^{\phi}\right)}{\partial \phi} \\
& =\frac{1}{r^{2}} \frac{\partial\left(r^{2} J^{r}\right)}{\partial r}+\frac{1}{\sin \theta} \frac{\partial\left(\sin \theta J^{\theta}\right)}{\partial \theta}+\frac{\partial J^{\phi}}{\partial \phi}
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Since $V=-r^{2} \sin \theta$ (see exercises).

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Now, note that
$J=J^{r} e_{r}+J^{\theta} e_{\theta}+J^{\phi} e_{\phi}=J^{r} \hat{e}_{r}+J^{\theta}\left(r \hat{e}_{\theta}\right)+J^{\phi}(r \sin \theta) \hat{e}_{\phi}=\hat{J}_{r} \hat{e}_{r}+\hat{J}_{\theta} \hat{e}_{\theta}+\hat{J}_{\phi} \hat{e}_{\phi}$

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Divergence

$$
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& =\frac{1}{r^{2}} \frac{\partial\left(r^{2} J^{r}\right)}{\partial r}+\frac{1}{\sin \theta} \frac{\partial\left(\sin \theta J^{\theta}\right)}{\partial \theta}+\frac{\partial J^{\phi}}{\partial \phi}
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Which agrees with the formula given in tables etc.

$$
\nabla \cdot J=\frac{1}{r^{2}} \frac{\partial\left(r^{2} \hat{J}_{r}\right)}{\partial r}+\frac{1}{r \sin \theta} \frac{\partial\left(\sin \theta \hat{J}_{\theta}\right)}{\partial \theta}+\frac{1}{r \sin \theta} \frac{\partial \hat{J}_{\phi}}{\partial \phi}
$$

## An Example: Spherical Polars in 3d cont...

Curl

$$
\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i}=
$$

$$
\left[\frac{\partial J_{r}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial r}\right]\left(e^{r} \wedge e^{\theta}\right)+\left[\frac{\partial J_{\phi}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial \phi}\right]\left(e^{\theta} \wedge e^{\phi}\right)+\left[\frac{\partial J_{r}}{\partial \phi}-\frac{\partial J_{\phi}}{\partial r}\right]\left(e^{\phi} \wedge e^{r}\right)
$$

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$$

Now look at the second component, noting that $e^{\theta} \wedge e^{\phi}=\frac{1}{r} \hat{e}_{\theta} \wedge \frac{1}{r \sin \theta} \hat{e}_{\phi}=\frac{1}{r^{2} \sin \theta} \hat{e}_{r} I$,

## An Example: Spherical Polars in 3d cont...

Curl

$$
\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i}=
$$

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$$
\left[\frac{\partial J_{\phi}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial \phi}\right]\left(e^{\theta} \wedge e^{\phi}\right)=\left[\frac{\partial\left(r \sin \theta \hat{J}_{\phi}\right)}{\partial \theta}-\frac{\partial\left(r \hat{J}_{\theta}\right)}{\partial \phi}\right] \frac{1}{r^{2} \sin \theta} \hat{e}_{r} I
$$

## An Example: Spherical Polars in 3d cont...

Curl

$$
\begin{gathered}
\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i}= \\
{\left[\frac{\partial J_{r}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial r}\right]\left(e^{r} \wedge e^{\theta}\right)+\left[\frac{\partial J_{\phi}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial \phi}\right]\left(e^{\theta} \wedge e^{\phi}\right)+\left[\frac{\partial J_{r}}{\partial \phi}-\frac{\partial J_{\phi}}{\partial r}\right]\left(e^{\phi} \wedge e^{r}\right)}
\end{gathered}
$$

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$$
\left[\frac{\partial J_{\phi}}{\partial \theta}-\frac{\partial J_{\theta}}{\partial \phi}\right]\left(e^{\theta} \wedge e^{\phi}\right)=\left[\frac{\partial\left(r \sin \theta \hat{J}_{\phi}\right)}{\partial \theta}-\frac{\partial\left(r \hat{J}_{\theta}\right)}{\partial \phi}\right] \frac{1}{r^{2} \sin \theta} \hat{e}_{r} I
$$

which can be written to agree with conventional tabulated form (though we have a bivector and not a vector):

## An Example: Spherical Polars in 3d cont...

Curl

$$
\nabla \wedge J=\left(\nabla J_{i}\right) \wedge e^{i}=
$$

$$
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$$
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$$

which can be written to agree with conventional tabulated form (though we have a bivector and not a vector):

$$
\frac{1}{r \sin \theta}\left[\frac{\partial\left(\sin \theta \hat{J}_{\phi}\right)}{\partial \theta}-\frac{\partial\left(\hat{J}_{\theta}\right)}{\partial \phi}\right]_{r} \hat{e}_{r} I
$$

## Exercises 5

(1) For 3d spherical polars, show that $V=-r^{2} \sin \theta$, where

$$
V I=\left(E^{n}\right)^{-1} \text { and } E^{n}=e^{r} \wedge e^{\theta} \wedge e^{\phi}
$$

(2) Show that the $e^{\phi} \wedge e^{r}$ component of $\nabla \wedge J$ can be written as:

$$
\frac{1}{r}\left[\frac{1}{\sin \theta} \frac{\partial\left(\hat{J}_{r}\right)}{\partial \phi}-\frac{\partial\left(r \hat{J}_{\phi}\right)}{\partial r}\right] \hat{e}_{\theta} I
$$

(3) Show that the $e^{r} \wedge e^{\theta}$ component on $\nabla \wedge J$ can be written as:

$$
\frac{1}{r}\left[\frac{1}{r} \frac{\partial\left(r \hat{J}_{\theta}\right)}{\partial r}-\frac{\partial\left(\hat{J}_{r}\right)}{\partial \theta}\right] \hat{e}_{\phi} I
$$

Check these against standard tabulated formulae.

## Connection with conventional Lamé Coefficients

Conventionally, sets of Lamé Coefficients are defined to be

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In our language we therefore have $h_{i}=\left|e_{i}\right|$, ie simply the moduli of the frame vectors defined by the coordinates.

All expressions of div, grad, curl etc in terms of the $h_{i} s$, can then be directly related to the expressions we derive in GA.

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- Shown how to construct Reciprocal Frames
- Using reciprocal frames, defined and motivated the form of the Vector Derivative
- Shown how to relate coordinates, frame vectors and reciprocal frame vectors in curvilinear coordinate systems.

